

# Global sensitivity analysis of lateral-torsional buckling resistance based on finite element simulations



Zdeněk Kala\*, Jan Valeš

Brno University of Technology, Faculty of Civil Engineering, Department of Structural Mechanics, Veverí Street 95, 602 00 Brno, Czech Republic

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## ABSTRACT

The article examines a hot-rolled steel I-beam subjected to lateral-torsional buckling (LTB) due to bending moment. The paper describes a non-linear finite element (FE) model and numerical approximation and simulation methods used for the global sensitivity analysis of the static resistance of a beam under major axis bending. The presented geometrically and materially non-linear FE model based on solid elements models in detail the LTB and the effects of initial imperfections on the ultimate limit state of a steel beam. Simulation runs of random imperfections are generated using the Latin Hypercube Sampling (LHS) method. Polynomial approximation of the model output helped minimise the number of runs of the non-linear finite element model. The approximation polynomial then facilitated the evaluation of sensitivity indices using a high number of simulation runs. The relationships between the slenderness and the first and second-order sensitivity indices are plotted in graphs. The graphs show the results of global sensitivity analyses of stochastic effects of initial imperfections and residual stress on the resistance of the investigated steel beam.

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## 1. Introduction

Lateral-torsional buckling (LTB) is a typical failure mode of unbraced I-beams subjected to bending about the major principal axis that both twist and deflect laterally. Lateral-torsional buckling resistance (LTB-R) including the effects of imperfections, residual stresses and plasticity effects may exhibit broad stochastic scattering due to the random influence of these effects [1]. Detrimental imperfections like initial curvature of the axis and residual stress may dramatically reduce the LTB-R. High sensitivity of LTB-R to the stochastic variability of initial imperfections increases the probability of failure of beams. Due to these reasons, it is important to study the effects of these inevitable initial imperfections and residual stresses on the LTB-R, which is the goal of the presented article.

The research of thin-walled beams has evolved from the classic Vlasov theory [2] to advanced non-linear finite element (FE) models; see e.g. [3–8]. In literature, the imperfect I-beam subjected to uniform bending moment is the only case for which the elastic closed-form analytical solution of lateral-torsional buckling resistance (LTB-E) exists [9]. The development of numerical

methods, such as FE methods, has lead to the continuous investigation of the available closed form solutions present in standard literature [9–13] and codes [14–16]. It has been clearly shown that such formulae are often conservative and in certain cases wrong [17].

The efficient use of computational models requires an understanding of key relationships between inputs and outputs. Analytical, especially numerical models of LTB failure modes are unfortunately often so complex that they prevent the understanding of the response of the model output to changes in model inputs based on intuition.

The dependence of model output on model inputs can be studied using stochastic sensitivity analysis (SA). The diversity of problems that are solved using mathematical modelling has recently lead to the development of several highly successful methods of sensitivity analysis [18–21]. Advanced methods of stochastic global sensitivity analysis are usually highly numerically demanding, especially in cases where it is necessary to evaluate the sensitivity indices using the Monte Carlo numerical simulation method. The computer time needed for the evaluation of one realization (one run) of the output random variable increases with increasing complexity of the computational model. Furthermore, the more complex the sensitivity analysis, the more realizations (numerical simulations) are needed to obtain statistically correct output.

\* Corresponding author.

E-mail addresses: [kala.z@fce.vutbr.cz](mailto:kala.z@fce.vutbr.cz) (Z. Kala), [vales.j@fce.vutbr.cz](mailto:vales.j@fce.vutbr.cz) (J. Valeš).

In engineering fields, scientists try to keep up to date with the current trends of modern research aimed at global sensitivity analysis (GSA) of model outputs [22]. GSA is necessary in situations where a priori information on the nature of the model is unavailable (model-free setting) or the model is considered to be non-additive [23]. In contrast, local SA (LSA) puts emphasis on the local (point) impact of the input factors on the model output [20]. LSA, which is also known as differential SA, belongs to the class of the one-factor-at-a-time methods. LSA has apparent advantages [21]. It is very straightforward compared to GSA, thus, making it easy to apply and interpret. In the case of a linear model, all information needed for SA is obtained from the first-order derivatives [24]. In the case of a non-linear but additive model, i.e. no interactions are present among factors, then derivatives of higher and cross order suffice to understand the model [24]. LSA cannot be used to describe higher-order interaction effects among factors. The outputs of non-linear and non-monotonous computational models such as LTB-R should be studied with GSA. One of the most effective methods of stochastic GSA is Sobol's sensitivity analysis SSA [25,26], which is based on the total decomposition of the variance of the output variable into terms with increasing dimensions.

In this article, SSA is used to map the influence of random input imperfections on the ultimate (inelastic) lateral-torsional buckling resistance (ultimate load carrying capacity) denoted as LTB-R. The term imperfection refers collectively here to both geometrical (initial axial twist and curvature, tolerance of dimensions of the cross section) and material (yield strength, residual stress, Young's modulus) imperfection [27]. The stochastic LTB-R was evaluated by Monte Carlo runs using Latin hypercube sampling (LHS) [28,29] on I-beam samples, which are randomly generated and subsequently analysed using the commercially available geometrically and materially non-linear FE program ANSYS [30].

A highly detailed model based on non-linear FE method allows detailed modelling of the effects of all imperfections on LTB-R. End-fork boundary conditions and uniform bending moment loading were considered, see Fig. 1. An elastic solution in the closed form [9], which was applied for the evaluation of LTB-E (lateral-torsional buckling resistance including the effects of imperfections, without residual stress and based on elastic and first yield analysis) [31,32], also exists for this type of boundary condition and loading. SSA was evaluated for LTB-R and LTB-E and the obtained results were compared. Due to the high complexity of the repeated calculations of LTB-R, SSA of the FE model was not performed, but rather the SSA of the polynomial approximation of the FE model.

## 2. FE model

Numerical simulations were performed using the Ansys software [30] as a process of running geometrically and materially non-linear model with regard to imperfections, such as initial geometric imperfections and residual stresses. The FE research was performed on a model of European hot-rolled steel I200 beam, which is statically loaded on both ends by bending moment  $M$ , see Fig. 1.

Due to potential problems that could occur whilst generating the FE mesh, the cross-sectional geometry of I200 was slightly simplified and fillets in the flange-to-web connection and at the ends of the flanges were neglected. The cross-sectional geometry of I200 is considered as biaxially symmetrical. The I200 cross-section is thus defined by five dimensions  $h$ ,  $b$ ,  $t_1$ ,  $t_2$  and  $\alpha$ , see Fig. 2. According to [33], the effect of the fillet was found to be negligible in LTB problems.

Commercially available software ANSYS was used for the FE analysis. Homogeneous structural solid element SOLID185 was used for the model. It is an 8-node element having three degrees of freedom at each node: translations in the nodal  $x$ ,  $y$ , and  $z$  direc-

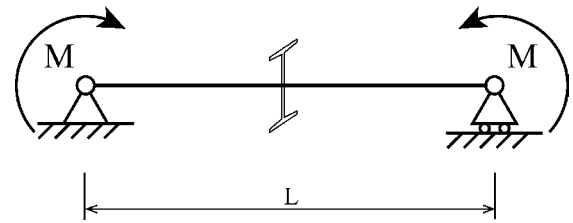


Fig. 1. Simply supported I-beam subjected to uniform bending moment.

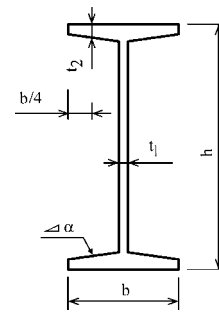


Fig. 2. The geometry of beam I200.

tions. It is suitable for 3D modelling of solid structures. It has plasticity, hyper-elasticity, large deflection and large strain capabilities, stress stiffening and creep. The enhanced strain formulation, which prevents shear locking in bending-dominated problems and volumetric locking in nearly incompressible cases, was considered. The element introduces nine internal degrees of freedom to handle shear locking, and four internal degrees of freedom to handle volumetric locking. All internal degrees of freedom are introduced automatically at the element level and condensed out during the solution phase of the analysis [30]. SOLID185 allows for prism, tetrahedral, and pyramid degenerations when used in irregular regions. However, this was not the case, and only hexahedron elements were created [30].

Even though it is customary to model LTB problems using shell elements instead of solid elements, certain undesirable effects, such as small material overlap in the flange-to-web connection, occur in shell models. The beam is thus a little bit more rigid and has a higher load-carrying capacity [3,34]. Another disadvantage of shell elements is the problematic modelling of the varying thickness of parts of the cross-section. In this case the flange thickness  $t_2$ . These undesirable effects were eliminated using solid elements SOLID185, however, at the cost of increasing demands on CPU time.

### 2.1. Mesh, boundary conditions and loads

Generally, the results obtained from the FE model are more accurate the finer the FE mesh. This, however, leads to an increase in the computational time. It is thus necessary to find a balance between the computational time and the desired accuracy. The mesh density was tested and the optimal setting of 20 elements over the entire height of the web, 10 elements over the entire width of the flanges and two elements over the web and flange thickness was chosen, see Fig. 3. It was tested that two elements over the thickness are sufficient and the final value of resistance is not influenced by the number of elements even in the plastic behaviour. The  $x$ -axis corresponds to the longitudinal axis of the beam; axes  $y$  and  $z$  lie in the plane of the cross-section. The origin of the coordinate system is placed in the centre of gravity of the cross-section. This also corresponds to the shear centre in a doubly

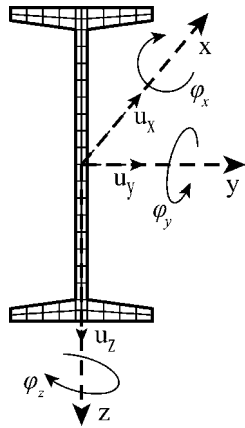


Fig. 3. The meshing of the cross-section I200.

symmetric cross-section. Translations  $u_x$ ,  $u_y$ ,  $u_z$  and rotations  $\phi_x$ ,  $\phi_y$ ,  $\phi_z$  are related to the global coordinate system. The number of elements per length was selected in a manner such that the maximum ratio of the longest to shortest side of the element (maxrat) was equal to 8. It may be noted that, according to [30], the aspect ratio for quadrilaterals should not exceed 20. Otherwise disturbance of analysis results may occur. This implies that the number of elements per beam length is not constant and the number of elements increases with increasing length of the beam. Control calculations were performed for a finer mesh with 4 elements over the thickness and maxrat = 4. The finer mesh yielded slightly more accurate results (difference of several tenths of a percent) for beams with low slenderness. However, the results from the present denser mesh are due to the small difference acceptable. For slender beams the results with the finer mesh were practically identical. For beams with  $\bar{\lambda}_{LT} = 2.0$  the accuracy of the FE model was also verified using the analytical solution in the closed form [31,32].

End-fork boundary conditions were considered for the model. These conditions make it possible to compare the elastic FE solution with the analytic solution [31,32] and verify the correctness of the model. Support conditions were created on both ends of the beam using three kinematic coupling constraints: two for the edges of the flanges and one for the axis of the web, see Fig. 4. Rotations  $\phi_z$  and  $\phi_y$  for the selected nodes of the web and flange respectively remain constant during loading, so that they are constrained to remain in a straight line. These boundary conditions were taken from the model presented in [3,34] and their implementation is crucial for the modelling of beams using solid or shell elements. Using these kinematic coupling constraints it is possible to overcome the occurrence of local extremes of stress at the end sections

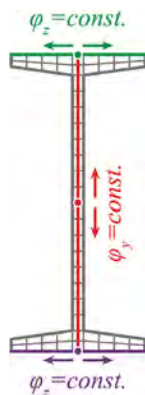


Fig. 4. Three kinematic coupling constraints.

during loading and certain convergence problems of the calculation. Alternatively, instead of using these constraints, it is possible in the end sections to model zones with gradually increasing values of kinematic hardening of the material diagram of steel towards the edges [35]. However, this approach is not as effective.

A symmetric model was considered. Longitudinal translation of a point in the centre of gravity was constrained to  $u_x = 0$  midspan of the beam. Boundary conditions of both ends of the beam are introduced into the centre of gravity and given as  $u_y = u_z = \phi_x = 0$ , see Fig. 5.

The beam is loaded on both ends by equal bending moments  $M$ . Moments are applied as surface loads in the form of pressure  $p$  [N/m<sup>2</sup>], which is defined by a gradient. The magnitude of the gradient is determined by the slope value (load per unit length) given as the ratio of the moment  $M$  to the moment of inertia  $I_y$ , and the distance from the  $z$ -axis, which represents the slope direction. The pressure acting on the element is given by the relationship  $p = M \cdot z / I_y$ . The pressure distribution is schematically shown in Fig. 5.

LTB-R was evaluated using the geometrically and materially non-linear FE method. The bending moment loading  $M$  was increased in steps until the value of  $M_{LTB-R}$ , where  $M_{LTB-R}$  is the maximum bending moment loading corresponding to the ultimate limit state.  $M_{LTB-R}$  is theoretically determined as the maximum moment  $M$  for which the determinant of the tangential stiffness matrix is equal to zero. Numerically,  $M_{LTB-R}$  is determined as the moment in the last loading step for which the determinant of the stiffness matrix is very close to zero whilst the calculation is still converging. The steps of the loading moment  $M_{LTB-R}$  are decreased so that the output value  $M_{LTB-R}$  is calculated with an accuracy of 0.3%.

## 2.2. Material model

The material stress–strain curves of structural steels are investigated and discussed in [34]. Steel of grade S235 was selected for the numerical study performed in this article. The hardening slope of  $E/10000$  was applied. The application of this hardening slope is based on a combination of numerical experience and recommendations of standards [36]. Annex C of [36] recommends the application of one of these four material behaviour models for plated FE analysis: (i) elastic–plastic without strain hardening, (ii) elastic–plastic with a nominal plateau slope of 1 MPa (or similarly small value), (iii) elastic–plastic with a strain hardening slope of  $E/100$ , (iv) true stress–strain curve modified from test results. The numerical instability of the FE computational model increases with

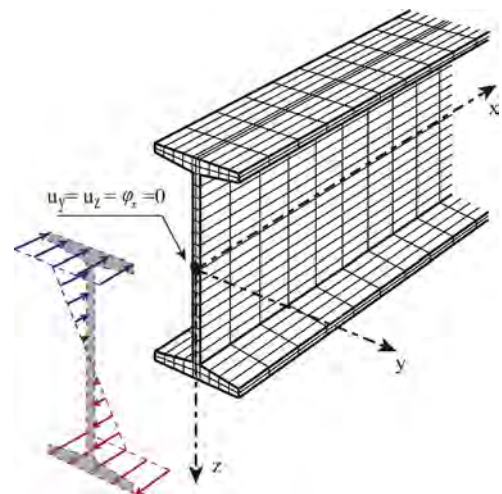


Fig. 5. The end-fork boundary conditions and equal bending moments.

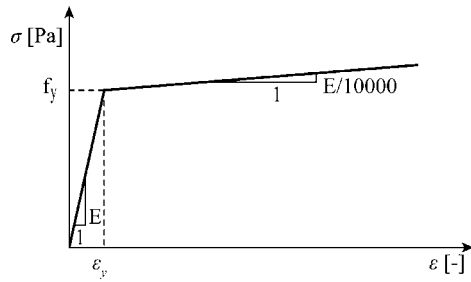


Fig. 6. Elastic-plastic stress-strain diagram of steel grade S235.

decreasing hardening slope, for e.g., the use of (i) lead to higher numerical instability than 1 MPa according to (ii). Furthermore, it was necessary to secure the numerical stability for each random realization (LHS simulation) of the input initial imperfections, which the application of (ii) did not fully ensure. The value of hardening slope, which is slightly greater than (ii) and provides sufficient numerical stability in each LHS simulation, was therefore numerically sought. The hardening slope of  $E/10,000$  was used to overcome numerical instability, see Fig. 6. For short beams the numerically obtained values of  $M_{LTB-R}$  based on  $E/10,000$  are virtually identical with the control results that were calculated using the plateau slope of 1 MPa (ii). It can be noted that the stress-strain relationship can be introduced in other ways, see e.g. [37,38,34].

### 2.3. Geometric imperfections

Consideration of geometrical out-of-straightness imperfection in the analysis of structural behaviour is very important. The out-of-straightness imperfection was obtained by scaling the first eigenvalue LTB mode shape of a perfectly straight uniformly bending beam. The “scale factor” of  $e_0$  was used to set the top flange maximum lateral imperfection at mid-span. The buckling mode was solved analytically as a pure elastic LTB problem; see e.g. [9,31,32]. The analytical solution very accurately matches the buckling mode obtained from the FE model. The advantage of the analytical solution is that the geometry in each LHS simulation can be quickly created automatically using the analytical formulas (1) and (2). Using the first eigenmode is very common in the investigation of the LTB phenomenon; see e.g. [4,39,6,40,41]. Both out-of-plane displacement  $v_0$  and torsional imperfection  $\varphi_0$  were

included. These initial imperfections are affine to the deformed shape and can be described using the following sine functions:

$$v_0 = a_{v0} \sin\left(\frac{\pi x}{L}\right), \quad \varphi_0 = a_{\varphi 0} \sin\left(\frac{\pi x}{L}\right) \quad (1)$$

where  $v_0$  is the curvature of the beam axis in the  $xy$ -plane and  $\varphi_0$  is the rotation of the cross-sections along the beam length, see Fig. 7.

The following relationships are valid for amplitudes  $a_{v0}$  and  $a_{\varphi 0}$  in a beam curved according to the first eigenmode

$$a_{v0} = \frac{|e_0|}{1 + \frac{h}{2} \frac{\pi^2 E I_z}{M_{cr} L^2}}, \quad a_{\varphi 0} = a_{v0} \frac{\pi^2 E I_z}{M_{cr} L^2} \quad (2)$$

where  $e_0$  is the amplitude at midspan relating to the centre of the top edge of the flange,  $h$  is the cross-section height,  $I_z$  is the second moment of area to the  $z$ -axis,  $L$  is the length of the beam,  $E$  is Young's modulus of elasticity and  $M_{cr}$  is the elastic critical moment at lateral beam buckling, see e.g. [9,31,32].

It may be noted that the procedure for the polynomial approximation described in Chapter 5 requires 400 random realizations of the geometry of the imperfect beam I200 for each length  $L$ . The random characteristics of geometry are flange thickness  $t_2$  and random amplitudes  $a_{v0}$  and  $a_{\varphi 0}$  (1), the other geometric characteristics are deterministic.  $a_{v0}$  and  $a_{\varphi 0}$  in (2) are randomly dependent on  $e_0$ ,  $E$ ,  $M_{cr}$ ,  $I_z$ .  $M_{cr}$  is randomly dependent on  $t_2$ ,  $E$  and  $I_z$  is randomly dependent on  $t_2$  [32]. Thus, the geometry of the I-beam is randomly dependent on three input random variables  $e_0$ ,  $t_2$  and  $E$ .

### 2.4. Residual stress

Residual stress can play a significant role in evaluating the structural behaviour and is strongly dependent on the manufacturing process. Parabolic or linear distribution of residual stress on the web and flanges may be applied to I-sections [42]. The commonly used residual stress pattern for hot-rolled I-profiles in FE modelling is the linear stress distribution [43], see Fig. 8.

The magnitude and distribution of residual stress in hot-rolled shapes depend on the type of cross-section, rolling temperature, cooling conditions, straightening procedures, and the material properties of the steel [13]. An overview of approaches for the modelling of residual stress can be found, for e.g. in [3,34,44]. Results of numerical studies found in [33,45] tend to relate the values of the residual stress to the material yield strength of steel grade S235. For the ratio  $h/b > 1.2$ , which corresponds to I profiles, the mean value of the residual stress may be considered as 30% of

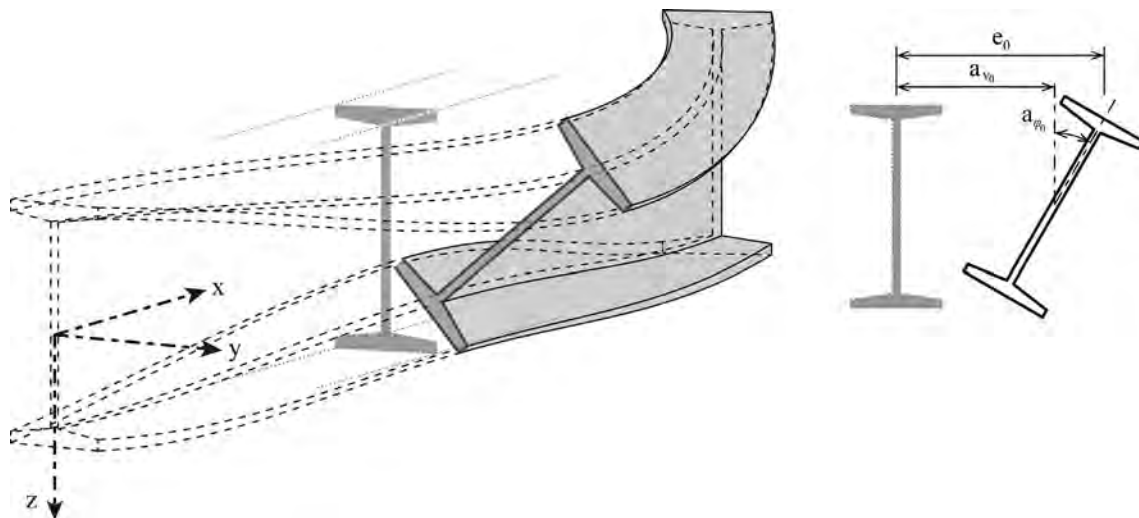


Fig. 7. Initial axis imperfection pattern based on the first mode of the elastic eigenvalue analysis.



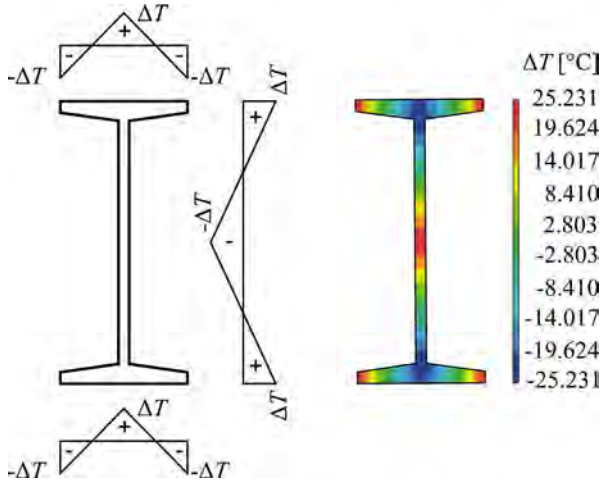


Fig. 8. Application of linear temperature changes  $\Delta T$  for the initiation of residual stress.

the yield strength  $f_y$ , for ratio  $h/b < 1.2$  (HE profiles) 50% of the yield strength is considered (for steel S235) [4,34]. Study [46] lists the values of residual stress at the flange edges of profiles IPE as  $0.3 \times 235$  MPa independent of the yield strength for steel S235, S355 and S460. Measurements of the residual stress in flanges of similar shapes, made from different steel grades, show that the distribution and size of the residual stress are very similar [13]. Results of experimental research have shown that the size of residual stress in hot-rolled profiles seems to be independent of the material yield strength for both mild and high-strength steel [34].

Residual stress of the hot-rolled I200 beam was introduced in the form of self equilibrating stress distribution through initial thermal load step. Change in temperature  $\Delta T$  at a point of the section is given by the formula

$$\Delta T = -\frac{\sigma_R}{E \cdot \alpha_t} \quad (3)$$

where  $\sigma_R$  is the residual stress in a given point and  $\alpha_t$  is the thermal expansion coefficient, considered as  $\alpha_t = 1.2 \times 10^{-5} \text{ [K}^{-1}\text{]}$ . Due to the overall configuration of the model the application of linear temperature change  $\Delta T$  yields the desired stress  $\sigma_R$  only approximately. The precise (required) value  $\sigma_R$  is obtained by a new setting of  $\Delta T$ , which is calculated by comparing the value attained (taken from the FE model) and the desired value of the residual stress when the dependency between  $\sigma_R$  and  $\Delta T$  is linear. Boundary conditions of both ends of the column with the described kinematic constraints are implemented after the introduction of the residual stress. The distribution of longitudinal stress  $\sigma_x$  initiated by change in temperature  $\Delta T$  is shown in Fig. 9.  $\sigma_R$  (blue)<sup>1</sup> is attained at the flange ends and in the centre of the web, whilst approximately  $0.8 \cdot \sigma_R$  (red) is obtained in the flange-to-web connection, see Fig. 9. Lower stress  $0.8 \cdot \sigma_R$  is due to the greater thickness of the flange in the flange-to-web connection, see Fig. 2. This difference does not occur in the case of beam models with rectangular flange shape (e.g. beams IPE, HEA, HEB), i.e. the absolute stress values from linear temperature change  $\pm \Delta T$  are the same at the end and the middle of the flange.

### 3. Sobol sensitivity analysis

Sobol's method has the capability to perform SA for non-linear, non-additive and non-monotonic models. Let  $Y = f(X_1, X_2, \dots, X_N)$  be

<sup>1</sup> For interpretation of color in Fig. 9, the reader is referred to the web version of this article.

a deterministic model, where  $Y$  is a scalar output and  $X_i$  are  $N$  statistically independent input variables. The variance of  $Y$  can be decomposed into the form:

$$V(Y) = \sum_i V_i + \sum_i \sum_{j>i} V_{ij} + \sum_i \sum_{j>i} \sum_{k>j} V_{ijk} + \dots + V_{123\dots N} \quad (4)$$

where  $V_i = V(E(Y|X_i))$  is the first order (main) partial variance and  $V_{ij} = V(E(Y|X_i, X_j)) - V_i - V_j$  is the second order partial variance etc.  $V_i$  can be explained as the average reduction of model output variance resulting from fixing  $X_i$ , this means that  $V_i$  measures the individual contribution of  $X_i$  to the total variance  $V(Y)$ . The second order partial variance  $V_{ij}$  measures the second order interaction contribution between  $X_i$  and  $X_j$  to  $V(Y)$ . The higher order partial variances for quantifying higher order interaction contributions can be expressed analogously. By normalising the partial variances  $V_i$ ,  $V_{ij}$ ,  $V_{ijk}$  etc., according to the total variance  $V(Y)$  we obtain Sobol sensitivity indices (5) of the first, second, third and higher orders.

$$S_i = \frac{V_i}{V(Y)}, S_{ij} = \frac{V_{ij}}{V(Y)}, S_{ijk} = \frac{V_{ijk}}{V(Y)}, \text{ etc.} \quad (5)$$

The total number of Sobol indices is  $2^n - 1$ . The sum of all Sobol indices must be equal to one.

$$\sum_i S_i + \sum_i \sum_{j>i} S_{ij} + \sum_i \sum_{j>i} \sum_{k>j} S_{ijk} + \dots + S_{1,2,3\dots N} = 1 \quad (6)$$

A more detailed description of Sobol's decomposition is in [25,26,22]. Examples of the application of SSA aimed at the stability and limit states of steel I-beams can be found in [32,47,48]. Although SSA is computationally demanding, it permits the exploration of all regions of the input space, accounting for interactions and non-linear responses.

### 4. Input random imperfections

The results of GSA [32] showed that LTB-E of an imperfect IPE-beam is most sensitive to the variability of four basic characteristics, namely yield strength (beams with low slenderness), Young's modulus (beams with high slenderness), flange thickness and the initial out-of-straightness of the beam axis. However, this conclusion was made on the basis of the analytical solution of LTB-E, in which the effects of residual stress were not modelled. Residual stress could be another significant imperfection if we were to study LTB-R.

The statistics of residual stress can be obtained from [49–51]. Three probability distribution models of residual stress evaluated from 103 measurements obtained from an extensive literature survey were defined in article [50]. Taking into account [50], we introduced two variants of the random variability of the residual stress. Residual stress at the flange edge is considered as input random variable  $rs$ . Gauss probability density function (pdf) with mean value  $\mu_{rs2} = 90$  MPa and coefficient of variation 0.2 (standard deviation  $\sigma_{rs2} = 18$  MPa) was considered in the variant  $rs_2$ . The residual stress is in the range of 10.5 MPa to 169.5 MPa if the LHS method is used to generate one hundred thousand numerical simulations. A four-parameter Hermite pdf [52] with  $\mu_{rs1} = 90$  MPa, coefficient of variation 0.3 ( $\sigma_{rs1} = 27$  MPa), skewness 0 and kurtosis 2.26 was considered in the variant  $rs_1$ . A residual stress range of 19.5–160.4 MPa was obtained from one hundred thousand numerical simulations generated using the LHS method, i.e. the observation interval is smaller than the Gauss pdf despite the larger standard deviation. Table 1 lists the comparison of statistical characteristics of variants  $rs_1$  and  $rs_2$  evaluated for 100 thousand LHS simulations. The comparison of both pdf types is presented in Fig. 10.

The statistical characteristics of yield strength of steel grade S235 were considered according to [53], where the results of

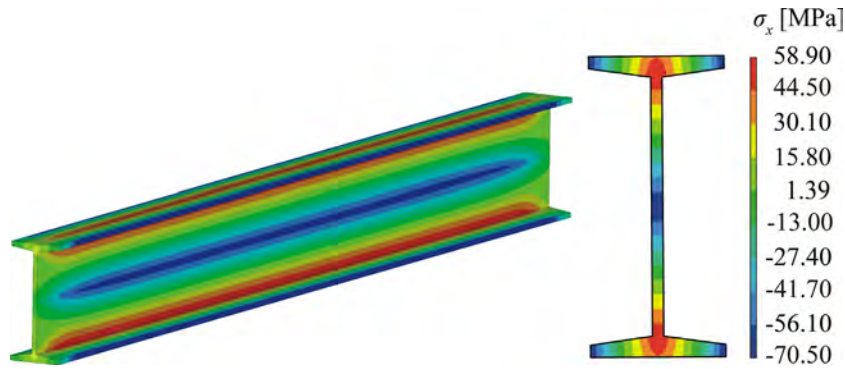


Fig. 9. Distribution of residual stress initiated by  $\Delta T$ .

**Table 1**  
Descriptive statistics of residual stress.

Variant names	$rs_1$	$rs_2$
Valid observations	100,000	100,000
Minimum	19.546	10.483
Maximum	160.45	169.52
Range	140.91	159.03
Median	90.000	90.000
Arithmetic mean	90.000	90.000
Geometric mean	85.452	88.094
Stand. deviation	27.259	18.000
Coef. of variation	0.3029	0.2000
Stand. skewness	-0.53E-08	-0.12E-12
Stand. kurtosis	2.2565	2.9995

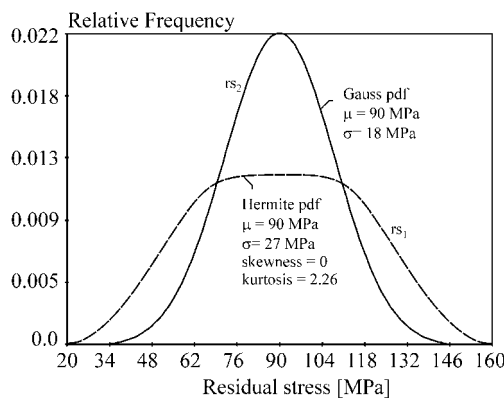


Fig. 10. The residual stress - Gauss and Hermite pdfs.

tensile tests of samples obtained from a third of flanges of profiles IPE 160 to IPE 220 were published. Note that the statistical evaluation of other independent experimental researches of the yield strength and other material characteristics of structural steel were published in [38,54]. The statistical characteristics of Young's modulus  $E$  were considered according to [55]. The statistical characteristics of flange thickness were considered according to [56].

The amplitude of initial imperfection  $e_0$  (2) was considered as a Gauss pdf with mean value of zero  $\mu_{e0} = 0$  (perfectly straight beam), which was taken from the JCSS Probabilistic Model Code [57,58]. Model Code [57] assumes that the pdf of  $e_0$  is symmetrical around zero and that small eccentricities are more likely than large ones, although the large ones are more dangerous. The standard deviation of  $e_0$  ( $\sigma_{e0}$ ) was derived from the assumption that 95 observations (random realizations) of this imperfection are found in the tolerance limits [32]. The out-of-straightness tolerance can be considered in accordance with long-term experience as  $0.1\% L$  [59,4,39]. This is consistent with codes of practice regulating the acceptability of production tolerances of steel structures. However,

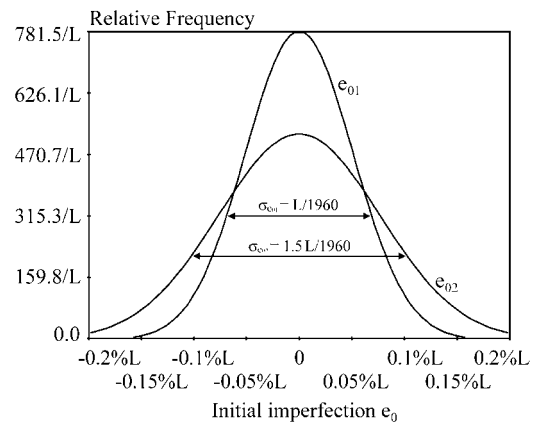


Fig. 11. The pdfs of initial imperfection  $e_0$ .

the straightness must comply with the requirements given in the European standard EN 10034:1993, which divides straightness tolerance  $q$  on length  $L$  into three groups according to height  $h$  [mm] for steel I and H sections:  $80 < h < 180$ ,  $q = 0.30\% L$ ;  $180 < h < 360$ ,  $q = 0.15\% L$ ;  $h > 360$ ,  $q = 0.1\% L$ . The tolerance of hot-rolled section I200 is  $q = 0.15\%$ . Therefore two variants  $0.1\% L$  and  $0.15\% L$  of tolerance were considered. In summary,  $\sigma_{e0}$  was considered in two variants as  $\sigma_{e01} = L/1960$  ( $0.1\% L$ ) [32,60] and  $\sigma_{e02} = 1.5L/1960$  ( $0.15\% L$ ) [31], see Fig. 11. It may be noted that the Gauss pdf is not the only possible choice for  $e_0$ . For example [61] introduces a log-normal pdf for  $e_0$  with  $\mu_{e0} = L/3000$ ,  $\sigma_{e0} = L/2551$  or  $\mu_{e0} = L/2000$ ,  $\sigma_{e0} = L/3785$ . The use of the Hermite pdf [52] with a shape similar to  $rs_1$  in Fig. 10, or a variant of the Hermite pdf with very small kurtosis that results in a “double-humped” and symmetrical around zero pdf may be discussed.

All input random geometric and material imperfections are listed in Table 2. Variables in Table 2 were considered as statistically independent. All other input geometric and material characteristics were considered using the nominal and characteristic (non-random) values of European standard I200 profile produced from steel of grade S235 [14], see Table 3.

## 5. Polynomial approximation of LTB-R

Sobol's method is much more computationally demanding in comparison to other global sensitivity analysis methods [21]. This is due to the fact that the series development of Eq. (6) has as many as  $2^n - 1$  terms. One possibility to evaluate all  $2^5 - 1 = 31$  Sobol indices is by approximating LTB-R using random realizations of the FE model output  $M_{LTB-R}$ . LTB-R is not calculated directly as  $M_{LTB-R}$ , but is represented by a polynomial (7) that has a quicker

response and which approximates all non-linear and interaction effects of non-linear FE model.

$$M_{LTB-R} \approx Y = \sum_{a=0}^2 \sum_{b=0}^2 \sum_{c=0}^2 \sum_{d=0}^2 \sum_{e=0}^2 c_{\alpha} \cdot X_1^a \cdot X_2^b \cdot X_3^c \cdot X_4^d \cdot X_5^e \quad (7)$$

The terms of the polynomial consist of the products comprising of  $n$ -combinations (for  $n = 0, 1, \dots, 5$ ) from a set of ten elements consisting of variables  $X_1, X_2, X_3, X_4, X_5$  and their squares, such that the exponent of each variable is at most two.

Approximation methods along with SA are not new and have been described, for e.g. with regards to the substitution of the limit state function using the Response Surface technique, which presents an acceptable compromise in terms of accuracy and computational effort, for the approximation of the limit state surface [62]. Polynomials suitable for the approximation of the limit state surface may contain non-linear terms, but usually do not retain the interaction effects of all higher orders. GSA and SSA generally require polynomials with more terms than approximations of the response surface type.

Polynomial (7) contains 5 input variables and has  $3^5 = 243$  terms with constants  $c_{\alpha}$ , where  $\alpha = 0, 1, \dots, 242$ . Polynomials with  $3^6 = 729$  terms for 6 input variables or polynomials with  $3^7 = 2187$  terms for 7 input variables, etc., can be expressed similarly. Approximation of the model output using polynomial (7) is advantageous for the evaluation of large numbers of Monte Carlo (LHS) runs, especially for computational models with a small number of input random variables that require large computing effort for the evaluation of a simulation (one run).

Constants  $c_{\alpha}$  were calculated using the least squares method. The LHS method was used to realize 400 simulations of random variables  $X_1, X_2, \dots, X_5$ , listed in Table 4, for which 400 runs of  $M_{LTB-R}$  were then simulated. The variances of random variables listed in Table 4 are higher than the variances of random variables listed in Table 2, which allows approximation (7) based on Table 4 to use high numbers of LHS simulation of the random variables from Table 2. The domain of the approximation is defined by the minimum and maximum values of Rectangular pdfs in Table 4, which were obtained as the minimum and maximum values from 500 thousand simulations of random imperfections from Table 2. This means that polynomial (7) can be evaluated for up to 500 thousand LHS simulations of initial imperfections from Table 2.

## 6. Sensitivity analysis of LTB-R

The goal of SSA is to study the effects of the variabilities of initial imperfections from Table 2 on LTB-R.  $V_i$  was evaluated using ten thousand realizations of  $E(Y|X_i)$ . One realization of  $E(Y|X_i)$  was performed by applying the approximation (7) and ten thousand LHS runs. The total variance  $V(Y)$  was evaluated using (7) and 500 thousand LHS runs. All Sobol higher-order indices were evaluated in a similar manner. Overall, all thirty-one Sobol indices were evaluated. The SSA procedure can be described as follows: the

non-dimensional slenderness of the beam  $\bar{\lambda}_{LT}$  is selected (e.g.  $\bar{\lambda}_{LT} = 0.9$ ) and the corresponding beam length is calculated  $L$  (e.g.  $L = 2.5$  m). The relationship between  $L$  and  $\bar{\lambda}_{LT}$  is listed in [14]. Lengths  $L$  are calculated according to the nominal parameters of an idealised cross-section, see Fig. 2. For I200 and  $\bar{\lambda}_{LT} \leq 1.7$  it approximately holds that  $L \approx 1.86\bar{\lambda}_{LT} + 0.41\bar{\lambda}_{LT}^2 + 0.71\bar{\lambda}_{LT}^3$ , see Table 3. The analytical formula is given in [63]. The length  $L$  is an input parameter of the formulae of standard deviations in Tables 2 and 4. Realization of random variables in Tables 2 and 4 are simulated using the LHS method. Due to the beam symmetry, LTB-R is the same for plus or minus realization  $e_0$ , therefore realizations  $e_0$  are considered using their absolute values in (2). The approximation polynomial (7) is compiled using Table 4. SSA of LTB-R is performed using (7) and Table 2. The variant of input imperfections with the probabilistic models corresponding closest to reality is  $t_2, f_y, E, e_{01}, rs_1$ , see Table 2. 31 Sobol indices are plotted in pie charts in Figs. 12–14, which clearly illustrate the changing effects of initial imperfections on LTB-R for three values of  $\bar{\lambda}_{LT}$ . The higher the value of Sobol's index, the higher the influence of the variability of initial imperfections on LTB-R. In order to obtain continuous curves of Sobol indices the step of 0.01 was considered for  $\bar{\lambda}_{LT}$ . First order sensitivity indices  $S_i$  are plotted in Fig. 15 and the crucial second order sensitivity indices  $S_{ij}$  are supplemented in Fig. 16. Similarly,  $S_i$  and  $S_{ij}$  for the other variants of initial imperfections from Table 2 are plotted; see Figs. 17–22. Sensitivity indices of the third and higher orders are relatively small. Thus, they are not plotted. It may be noted that the results in Figs. 17–22 do not depict the effects of variables  $h, b, t_1$  because their random variabilities were neglected. This is in accordance with [32,35,64], where the small effect of the variability of variables  $h, b, t_1$  on LTB-E and LTB-R was proven.

**Table 3**  
Nominal geometric and material characteristics of I200.

Symbol	Characteristic	Value
$h$	Sectional height	200 mm
$b$	Sectional width	90 mm
$t_1$	Web thickness	7.5 mm
$\alpha$	Inclination angle	14%
$\nu$	Poisson's ration	0.3
$G$	Shear modulus	$E/(2(1 + 0.3))$
$L$	Beam length	$\approx 1.86\bar{\lambda}_{LT} + 0.41\bar{\lambda}_{LT}^2 + 0.71\bar{\lambda}_{LT}^3$

**Table 4**  
Artificial random variables for approximation.

Symbol	Characteristic	Density	Minimum	Maximum
$X_1$	Flange thickness	Rectangular	8.83 mm	13.77 mm
$X_2$	Yield strength	Rectangular	217.43 MPa	377.17 MPa
$X_3$	Modulus of elasticity	Rectangular	162.46 GPa	257.54 GPa
$X_4$	Initial imperfection	Rectangular	−4.76 L/1960	4.76 L/1960
$X_5$	Residual stress	Rectangular	4 MPa	176 MPa

**Table 2**  
Input random imperfections.

Symbol	Characteristic	Density	Mean $\mu$	Standard deviation	Variant
$t_2$	Flange thickness	Gauss	11.3 mm	0.518 mm	$t_2$
$f_y$	Yield strength	Gauss	297.3 MPa	16.8 MPa	$f_y$
$E$	Modulus of elasticity	Gauss	210 GPa	10 GPa	$E$
$e_0$	Initial imperfection	Gauss	0	L/1960	$e_{01}$
		Gauss	0	1.5L/1960	$e_{02}$
$rs$	Residual stress	Hermite	90 MPa	27 MPa	$rs_1$
		Gauss	90 MPa	18 MPa	$rs_2$

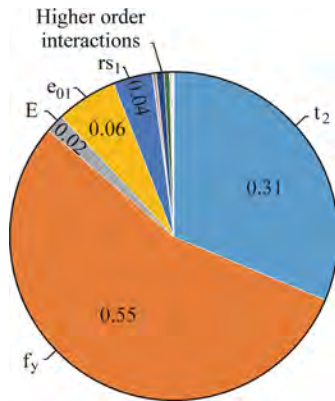


Fig. 12. SSA of LTB-R for variant  $t_2, f_y, E, e_{01}, rs_1$  and  $\bar{\lambda}_{LT} = 0.6$ .

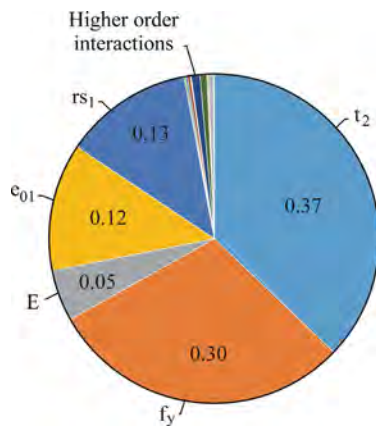


Fig. 13. SSA of LTB-R for variant  $t_2, f_y, E, e_{01}, rs_1$  and  $\bar{\lambda}_{LT} = 0.8$ .

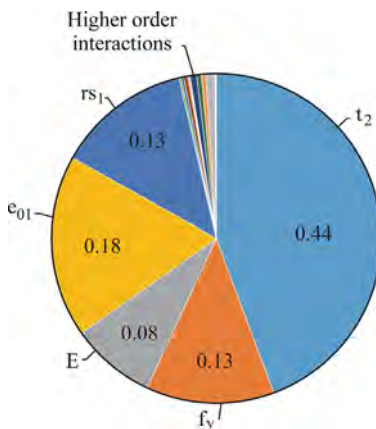


Fig. 14. SSA of LTB-R for variant  $t_2, f_y, E, e_{01}, rs_1$  and  $\bar{\lambda}_{LT} = 1.0$ .

The thick green curves are the SSA outputs of LTB-E. LTB-E is denoted in [32] as  $M_R$  (elastic load-carrying capacity) and is derived in the closed form. Let us note that the correctness of the analysis of  $M_R$  was verified using the geometrically non-linear FE model [63]. The analytical elastic solution  $M_R$  does not take into account the effect of residual stress. The input variables are otherwise the same for SSA, see Table 2.

Regarding the residual stress, the introduction of a high standard deviation  $\sigma_{rs1} = 27$  MPa did not lead (as expected) to a high

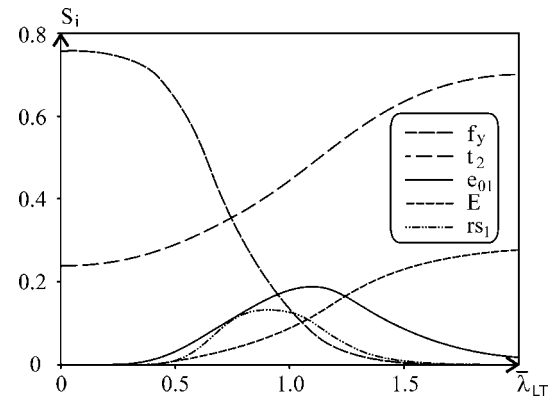


Fig. 15. First order indices  $S_i$  for variant  $t_2, f_y, E, e_{01}, rs_1$ .

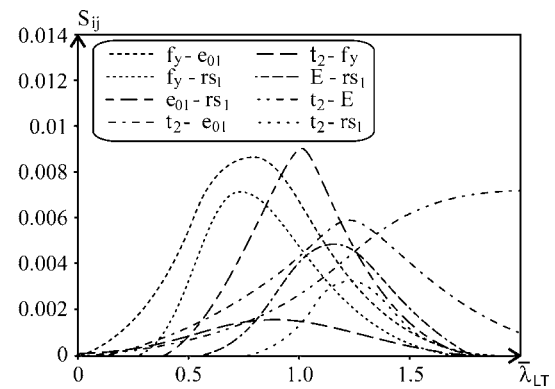


Fig. 16. Crucial second order indices  $S_{ij}$  for variant  $t_2, f_y, E, e_{01}, rs_1$ .

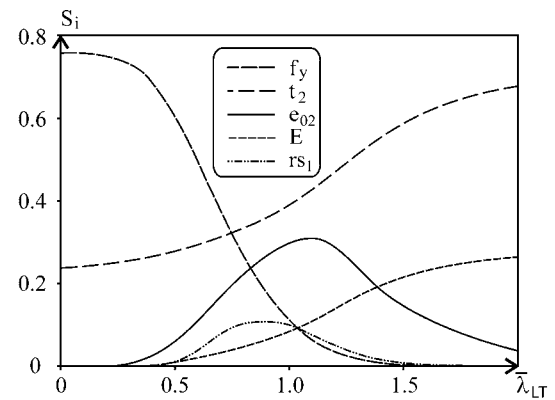


Fig. 17. First order indices  $S_i$  for variant  $t_2, f_y, E, e_{02}, rs_1$ .

value of sensitivity index  $S_{rs1}$ , see Figs. 15 and 17. Sensitivity of LTB-R to  $rs_1$  is relatively small. It is essential to properly understand this observation. SSA evaluates the effect of an imperfection while all other imperfections are varying as well. Whatever the strength interaction in the model,  $S_{rs1}$  indicates by how much one could, on an average, reduce the variance of LTB-R if  $rs_1$  is fixed. Let us perform a comparison with deterministic study [65] investigating the effect of the inclusion or exclusion of the nominal value of residual stress on LTB-R of beam I200 with nominal imperfections. The residual stress reduced LTB-R by 6% for  $\bar{\lambda}_{LT} = 1.2$ , the effect was smaller for other values of slenderness.



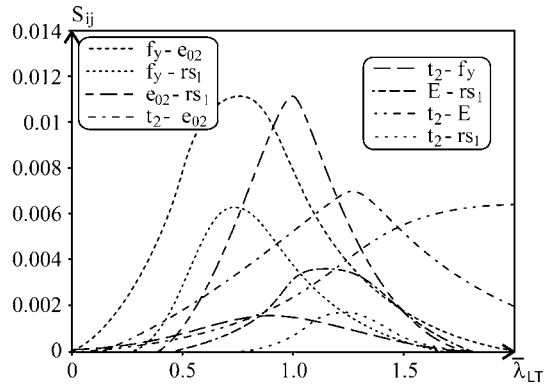


Fig. 18. Crucial second order indices  $S_{ij}$  for variant  $t_2, f_y, E, e_{02}, rs_1$ .

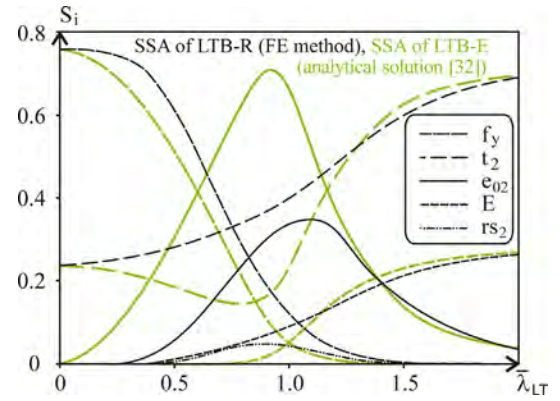


Fig. 21. First order indices  $S_i$  for variant  $t_2, f_y, E, e_{02}, rs_2$ .

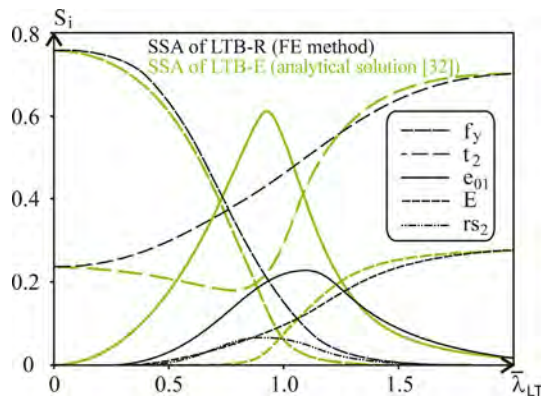


Fig. 19. First order indices  $S_i$  for variant  $t_2, f_y, E, e_{01}, rs_2$ .

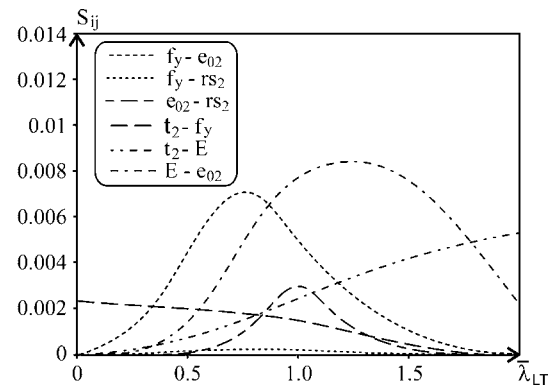


Fig. 22. Crucial second order indices  $S_{ij}$  for variant  $t_2, f_y, E, e_{02}, rs_2$ .

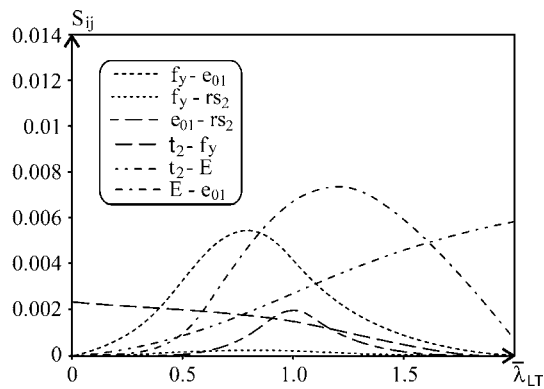


Fig. 20. Crucial second order indices  $S_{ij}$  for variant  $t_2, f_y, E, e_{01}, rs_2$ .

This study belongs to the LSA category, because the value of the residual stress changes while all other imperfections are fixed. The effect of residual stress can generally be different if the other imperfections are set to different values. Using a two-valued change of another imperfection (e.g.  $e_0$  as 0 and  $L/1000$ ) can have a far greater effect on LTB-R than a two-valued change in residual stress. Changing the values of only some imperfections while keeping other imperfections fixed (no-change) can lead to the misinterpretation of their joint influence. Selection of the base case is very important in LSA. The SSA presented above provides much more information because it quantifies the overall contribution of each random input imperfection to the uncertainty of LTB-R.

## 7. Observations

The paper describes the detailed non-linear FE model and advanced numerical methods that were used for the in-depth sensitivity analysis of LTB-R. Table 2 lists the input random initial imperfections of the beam I200, for which SSA was performed. The results of SSA are useful for the comparison of the relative importance of the five initial imperfections, which are the model inputs of the LTB-R analysis. Unless stated otherwise, the conclusions of the SSA of LTB-R pertain to all solution variants; see Figs. 15–22.

- The influence of the residual stress  $rs_1$  on LTB-R is relatively small despite the fact that random variable  $rs_1$  was introduced with large variance. The same conclusion is true for random variable  $rs_2$ , which has lower (but also quite large) variance. The curves of sensitivity indices  $S_{rs1}$  and  $S_{rs2}$  reach their maximum for approximately  $\bar{\lambda}_{LT} = 0.9$ .
- Sensitivity index of yield strength  $S_{f_y}$  has a plateau near  $\bar{\lambda}_{LT} = 0$  and then decreases with increasing  $\bar{\lambda}_{LT}$ . For approximately  $\bar{\lambda}_{LT} > 1.7$   $S_{f_y} = 0$ . This means that the variance of the yield strength can be fixed at any value of the domain without any effect on the variance of LTB-R.
- The influence of imperfection  $e_0$  on LTB-R is the second dominant for  $\bar{\lambda}_{LT} \approx 1.1$ , especially for the variant  $e_{02}$ , which has 50% higher standard deviation than variant  $e_{01}$ . The curves of sensitivity indices  $S_{e01}$  and  $S_{e02}$  attain their maximum approximately for  $\bar{\lambda}_{LT} = 1.1$ . Sensitivity indices  $S_{e01}$  and  $S_{e02}$  are clearly higher than indices  $S_{rs1}$  and  $S_{rs2}$ .

- Sensitivity index  $S_{t_2}$  of the flange thickness  $t_2$  increases with increasing  $\bar{\lambda}_{LT}$ .  $t_2$  is the absolutely dominant imperfection for  $\bar{\lambda}_{LT} > 0.75$ . This means that reducing the variance of  $t_2$  has a relatively significant effect on the reduction of the variance of LTB-R. Reducing the variance of  $t_2$  can be achieved by improving the production processes in metallurgical production.
- Sensitivity index of Young's modulus  $S_E$  increases with increasing  $\bar{\lambda}_{LT}$ . In contrast to  $t_2$ , the variance of  $E$  cannot be significantly influenced through technological processes in metallurgical production.

Second order sensitivity indices  $S_{ij}$  measure the second order interaction contribution between  $X_i$  and  $X_j$  to the variance of LTB-R. The values of  $S_{ij}$  shown in Figs. 16, 18, 20 and 22 are small, but provide interesting information on the non-linear response of an imperfect beam in limit state. Sensitivity to the yield strength  $f_y$  and the involvement of  $f_y$  in interactions determine, among other things, the sensitivity to plastic yielding. Plastic yielding does not occur when  $S_{f_y}$  and fifteen additional higher-order indices involved in interactions with  $f_y$  are equal to zero. This is approximately fulfilled for  $\bar{\lambda}_{LT} > 2.0$ , otherwise at least one of the indices  $S_{f_y}$ ,  $S_{t_2, f_y}$ ,  $S_{f_y, e_0}$ ,  $S_{f_y, r_s}$  always has a non-zero value. The relatively small values of all 26 higher-order sensitivity indices show that the non-linear model can be considered as approximately additive. This is favourable information for the selection of the approximation strategies that use second-order polynomials neglecting interaction effects. This conclusion applies to an individual beam; beams integrated into structural systems may be influenced by interactions with surrounding structures [66]. The thick green curves (Figs. 19 and 21) show the results of SSA of LTB-E, which were calculated using the elastic analytic solution in the closed form [31,32]. The thick green curves are very different in comparison with the black curves (SSA of LTB-R). The analytic elastic solution greatly exaggerates the importance of initial imperfections  $e_0$  and on the contrary under-values the importance of  $t_2$  for the intermediate slenderness. This difference can be explained by the varied involvement of the flanges in the LTB-E and LTB-R solution. The thick green and black curves correspond for  $\bar{\lambda}_{LT} = 0$  and  $\bar{\lambda}_{LT} = 2.0$ . The thick green curves of  $S_{e_01}$  and  $S_{e_02}$  are maximum for  $\bar{\lambda}_{LT} = 0.93$  coinciding well with the maximum of  $S_{r_s1}$  and  $S_{r_s2}$ , which occurs in  $\bar{\lambda}_{LT} = 0.9$ . The thick green curves of  $S_{e_01}$  and  $S_{e_02}$  are also very similar in shape to the black curves of  $S_{r_s1}$  and  $S_{r_s2}$ . This is consistent with the frequently applied assumption that LTB-R taking into account the effect of residual stress can be approximately calculated as LTB-E of a beam with increased amplitude of initial imperfection  $e_0$ , sometimes referred to as the equivalent geometric imperfection.

## 8. Conclusion

GSA is an important part of the analysis of the reliability of steel structures and their limit states. SSA answers the question of which imperfections are dominant and should thus be paid increased attention during the preparation of the input values of stochastic models. Despite ongoing experimental research probabilistic models of initial curvature of the beam axis and residual stress are still under discussion. These imperfections influence the LTB-R, especially for I-beams with intermediate slenderness, which are commonly found in steel structures.

SSA showed that the random variability of residual stress  $r_s$  has a smaller influence on LTB-R than the flange thickness  $t_2$ , axis curvature  $e_0$  and yield strength  $f_y$ . Thus the statistical characteristics of residual stress need not be detected with such accuracy as  $t_2$ ,  $e_0$  and  $f_y$ . This conclusion is optimistic, because as we know, experimental research of residual stress is very difficult, time-consuming and an inefficient task with limited accuracy [44].

SSA showed that the random variability of  $e_0$  has in some cases the second dominant effect on LTB-R. Thus the statistical characteristics and pdf type of  $e_0$  must be determined with great accuracy, otherwise the results of probabilistic reliability analyses may be burdened with high error. The elastic solution LTB-E [31,32] overstates the influence of  $e_0$ , understates the importance of  $t_2$  and does not take into account the effect of residual stress. For certain  $\bar{\lambda}_{LT}$  the sensitivity of LTB-E to  $e_0$  is more than twice higher than the sensitivity of LTB-R to  $e_0$ . Comparison of results of sensitivity analyses LTB-R and LTB-E can be useful for the probabilistic verification of stability design criteria of steel structures.

The results of SSA of LTB-R were obtained using the polynomial approximation of a FE model, which is a highly effective approach especially for numerically demanding models with a small number of input random variables. This approach made it possible to handle detailed numerical observation of the effects of imperfections on LTB-R, which could not otherwise be achieved.

## Acknowledgements

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